

Escape fractions from unattenuated Ly α emitters around luminous $z > 6$ quasars

MINGHAO YUE,¹ ANNA-CHRISTINA EILERS,¹ JORRYT MATTHEE,² ROHAN P. NAIDU,^{1,*} RONGMON BORDOLOI,³
FREDERICK B. DAVIES,⁴ JOSEPH F. HENNAWI,^{5,6} DAICHI KASHINO,⁷ RUARI MACKENZIE,⁸ AND ROBERT A. SIMCOE¹

¹MIT Kavli Institute for Astrophysics and Space Research, 77 Massachusetts Ave., Cambridge, MA 02139, USA

²Institute of Science and Technology Austria (ISTA), Am Campus 1, 3400 Klosterneuburg, Austria

³Department of Physics, North Carolina State University, Raleigh, 27695, North Carolina, USA

⁴Max Planck Institut für Astronomie, Königstuhl 17, D-69117 Heidelberg, Germany

⁵Department of Physics, University of California, Santa Barbara, CA 93106-9530, USA

⁶Leiden Observatory, Leiden University, P.O. Box 9513, 2300 RA Leiden, The Netherlands

⁷National Astronomical Observatory of Japan, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan

⁸Institute of Physics, Laboratory of Astrophysics, Ecole Polytechnique Fédérale de Lausanne (EPFL),

and

Observatoire de Sauverny, 1290 Versoix, Switzerland.

ABSTRACT

Ionized proximity zones around luminous quasars provide a unique laboratory to characterize the Ly α emission lines from $z > 6$ galaxies without significant attenuation from the intergalactic medium (IGM). However, Ly α line measurements for galaxies within high-redshift quasars' proximity zones have been rare so far. Here we present deep spectroscopic observations obtained with the NIRSpec/MSA instrument on the James Webb Space Telescope (JWST) of galaxies in two $z > 6$ quasar fields. We measure the Ly α line fluxes for 50 galaxies at $6 < z < 7$ with UV absolute magnitude $M_{\text{UV}} < -19$ (median $M_{\text{UV}} = -19.97$), among which 15 are located near the luminous quasars, i.e. within $\Delta v < 2500 \text{ km s}^{-1}$. We find that galaxies near the quasars show significant flux bluewards of the systemic Ly α wavelength, and have **higher Ly α equivalent width** compared to galaxies at similar redshifts that are not located within the quasars' environment. Our result indicates little or no redshift evolution for the Ly α -emitter fraction from $z \sim 6.4$ to $z \sim 5$. Leveraging the low IGM opacity in the quasars' vicinity, we evaluate the Ly α escape fraction ($f_{\text{esc}}^{\text{Ly}\alpha}$) of high-redshift galaxies. Our analysis suggests that galaxies at $\langle z \rangle \approx 6.4$ have an average $f_{\text{esc}}^{\text{Ly}\alpha} = 0.14 \pm 0.04$. This value is consistent with reionization models where the Lyman continuum escape fraction is low ($f_{\text{esc}}^{\text{LyC}} \lesssim 0.1$) for luminous galaxies, and where the most luminous galaxies have only a minor contribution to the total ionizing photon budget.

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1. INTRODUCTION

The Epoch of Reionization (EoR) represents the last major phase transition of the universe, where the neutral hydrogen fraction of the intergalactic medium (IGM) changes from ~ 1 to ~ 0 between redshift $z \sim 20$ and $z \sim 5.5$ (e.g., S. E. I. Bosman et al. 2022; B. E. Robertson 2022). The overall history of reionization has been constrained by various observations, including the Cosmic Microwave Background measurements (e.g., Planck Collaboration et al. 2020), quasar transmission fluxes and damping wings (e.g., X. Fan et al. 2006; J. Yang et al. 2020; X. Jin et al. 2023; D. Āurovčíková et al.

2024), and Ly α lines from high-redshift galaxies (e.g., C. A. Mason et al. 2018, 2019; I. G. B. Wold et al. 2022; M. Tang et al. 2024a; Y. Kageura et al. 2025; H. Umeda et al. 2025). Nonetheless, it is still under debate which type of sources dominates the ionizing photon production. Specifically, several studies have suggested that luminous galaxies make a major contribution to reionization (e.g., M. Sharma et al. 2016; R. P. Naidu et al. 2020; J. Matthee et al. 2022), while others found that the ionizing photon production is dominated by faint galaxies (e.g., R. J. Bouwens et al. 2012; S. L. Finkelstein et al. 2019; H. Atek et al. 2024).

One reason for this debate is the highly uncertain Lyman continuum (LyC) escape fraction ($f_{\text{esc}}^{\text{LyC}}$) for the EoR galaxies. Due to the high IGM opacity to LyC photons at $z \gtrsim 6$, it is impossible to directly measure

Corresponding author: Minghao Yue

* NHFP Hubble Fellow

the leaking LyC flux from EoR galaxies. Several indirect methods have been employed to infer the escape fraction of galaxies at $z \gtrsim 6$, including observations for low-redshift counterparts of $z > 6$ galaxies (e.g., C. C. Steidel et al. 2018; R. P. Naidu et al. 2018; Y. I. Izotov et al. 2018), computing the ionizing photon budget that is needed to produce the observed neutral fraction evolution (e.g., S. L. Finkelstein et al. 2019; R. P. Naidu et al. 2020; B. E. Robertson 2022), and adopting predictions from cosmological simulations (e.g., X. Ma et al. 2015, 2016; J. Rosdahl et al. 2022). Nevertheless, the inferred LyC escape fraction shows large discrepancies, ranging from $f_{\text{esc}}^{\text{LyC}} \lesssim 5\%$ (e.g., J. Chisholm et al. 2022; S. Mascia et al. 2023; C. Simmonds et al. 2024; C. Papovich et al. 2025) to $f_{\text{esc}}^{\text{LyC}} \gtrsim 20\%$ (e.g., R. P. Naidu et al. 2020; A. E. Jaskot et al. 2024).

At $4 \lesssim z \lesssim 5$, where the IGM is mostly opaque to LyC photons but has high transmission to Ly α lines from galaxies (e.g., R. A. Meyer et al. 2025), the LyC escape fraction of galaxies can be estimated using the escape fraction of the Ly α emission line ($f_{\text{esc}}^{\text{Ly}\alpha}$; e.g., Y. I. Izotov et al. 2022; R. Begley et al. 2024). However, this method is not applicable to most galaxies in the EoR, as the Ly α line is heavily attenuated by the neutral hydrogen in the $z > 6$ IGM. The only viable approach to measure the unattenuated Ly α lines from EoR galaxies is to target galaxies that reside in ionized bubbles, where the IGM is transparent to Ly α photons (e.g., J. Matthee et al. 2018; T.-Y. Lu et al. 2024; L. Whitler et al. 2024; A. Torralba et al. 2024; Z. Chen et al. 2025). Such ionized bubbles are often produced by galaxy overdensities, and finding these ionized bubbles requires deep spectroscopy over wide fields (e.g., I. Nikolić et al. 2025).

In this context, proximity zones of high-redshift quasars offer a unique opportunity to investigate the unattenuated Ly α properties of EoR galaxies (e.g., S. E. I. Bosman et al. 2020; K. Protušová et al. 2024). Quasars at $z > 6$ are expected to produce large ionized bubbles around them, as implied by radiative transfer simulations (e.g., F. B. Davies 2020). Deep spectroscopy for galaxies residing in the proximity of high-redshift quasars will put direct constraints on the Ly α line properties of these galaxies.

In this work, we present *James Webb Space Telescope* (*JWST*) NIRSpect/MSA spectroscopy for galaxies near two $z > 6$ quasars, i.e. J0100+2802 ($z = 6.327$) and J1148+5251 ($z = 6.42$). We analyze the Ly α line properties of these galaxies, and try to answer the following questions: (1) For galaxies around high-redshift quasars, can we detect their Ly α emission lines without significant IGM attenuation? (2) How can we use galaxies

near $z > 6$ quasars to constrain the Ly α escape fraction of EoR galaxies?

This paper is organized as follows. Section 2 describes the NIRC*am* and NIRS*pec* observations used in this study. In Section 3, we describe how we measure the Ly α line properties of the targeted galaxies, and present evidence of less-attenuated Ly α lines from galaxies near the luminous quasars. In Section 4, we evaluate the escape fraction of galaxies near quasars utilizing the high IGM transmission around quasars, and discuss the implications for reionization models. We make conclusions in Section 5. Throughout this paper, we adopt a flat Λ CDM cosmology with $H_0 = 70 \text{ km s}^{-1}$ and $\Omega_M = 0.3$, and use the AB magnitude system (J. B. Oke & J. E. Gunn 1983).

2. DATA

2.1. NIRC*am* Observations

This work focuses on galaxies observed in two high-redshift quasar fields, namely around the quasars J0100+2802 ($z = 6.327$) and J1148+5251 ($z = 6.42$). These quasar fields were observed as part of the *Emission-line galaxies and Intergalactic Gas in the Epoch of Reionization* (EIGER) program (GO-1243, PI: Lilly), which delivers NIRC*am* F115W, F200W, and F356W imaging and F356W grism spectroscopy. Figure 1 shows the NIRC*am* images of the two quasar fields. The images have sizes of $3' \times 6'$, which corresponds to $\approx 1 \times 2$ proper Mpc at the quasars' redshifts. We refer the readers to D. Kashino et al. (2023) and J. Matthee et al. (2023) for details on the data reduction of the EIGER observations. The typical 5σ depths of the F115W, F200W, and F356W imaging are 27.8, 28.3, and 28.1 magnitudes, respectively. The sensitivity of the F356W grism spectra depends on the wavelength and the source positions, with the best sensitivity reaching $0.6 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2}$ at 3.8 micron.

Using the F356W grism spectra, J. Matthee et al. (2023) and D. Kashino et al. (2025) identified [O III]-emitters in the two quasar fields and performed spectral energy distribution (SED) fitting for these galaxies. In this work, we use the [O III] redshifts as the systemic redshifts (z_{sys}) of galaxies. We also estimate the absolute magnitude at rest-frame 1500\AA (M_{UV}) of all the galaxies by fitting a power law SED to the F115W and F200W magnitudes. These galaxy properties can be found in Table 1.

2.2. NIRS*pec*/MSA Spectroscopy

We obtain the NIRS*pec*/MSA G140M/F070LP spectroscopy for the two quasar fields as part of the *Mapping Super-luminous Quasars' Extended Radiative Emission*

(MASQUERADE) project (GO-3117 and GO-4713, PI: Eilers). The spectra cover wavelengths from $8000\text{\AA} \lesssim \lambda_{\text{obs}} \lesssim 13000\text{\AA}$, with a wavelength-dependent spectral resolution of $R \approx 600 - 1000$. At the quasars' Ly α wavelengths ($\lambda_{\text{obs}} \approx 9000\text{\AA}$), the spectral resolution is $R \approx 600$. Each quasar field is covered by two MSA pointings, and we adopt the 3-Shutter dither pattern for the observations. The total on-source exposure time on each pointing is 7.7 hours.

The primary targets of the MSA observation are [O III]-emitters at $z_{\text{sys}} > 5.5$ in the quasar fields. In this study, we focus on galaxies at $z_{\text{sys}} > 6$ and with $M_{\text{UV}} < -19$, as the MSA observations are not sufficiently deep to detect the continuum of fainter sources. The circles in Figure 1 show the location of the galaxies we analyze in this work, and the galaxy coordinates can be found in Table 1.

We reduce the spectra by running the `Detector1Pipeline` using `rwst`⁹ version 1.17.1 and CRDS `rwst_1322.pmap` to get the rate files, then using `msaexp`¹⁰ version 0.9.5.dev8+ge2b237b (G. Brammer 2023) to extract the 2D spectra. As the final step, we use `PypeIt` (J. X. Prochaska et al. 2020)¹¹ version 1.17.1 to coadd the 2D spectra for individual objects and extract their 1D spectra.

3. UNATTENUATED Ly α LINES FROM GALAXIES NEAR QUASARS

Luminous quasars at $z > 6$ are expected to produce large ionized bubbles around them, where the IGM has low opacity to Ly α photons. In this section, we test whether we can measure the Ly α lines from galaxies near quasars without significant IGM attenuation. We first fit the galaxy spectra to measure their Ly α fluxes, then compare the Ly α profiles and equivalent widths of galaxies near quasars to galaxies outside the quasars' ionized bubbles.

3.1. Fitting the Ly α Lines

We fit the Ly α lines of galaxies targeted by the NIR-Spec/MSA observation. We exclude five objects that have their Ly α lines falling in the detector gap, and three objects that are severely contaminated by nearby bright sources. For each remaining object, we fit the spectral window $1200\text{\AA} < \lambda_{\text{rest}} < 1500\text{\AA}$, where $\lambda_{\text{rest}} = \lambda_{\text{obs}}/(1 + z_{\text{sys}})$ is the rest-frame wavelength. We use a power-law continuum plus a Gaussian profile to model the spectral region around the Ly α emission, where the

continuum flux at $\lambda_{\text{rest}} < \lambda_{\text{Ly}\alpha}$ is set to be zero. This model contains five free parameters, namely an amplitude (A) and a slope (β) of each galaxy's continuum, the Ly α redshift ($z_{\text{Ly}\alpha}$), the Ly α flux ($F_{\text{Ly}\alpha}$), and the line width ($\sigma_{\text{Ly}\alpha}$). We run a nested sampling algorithm `dynesty` to determine the posterior probability distribution for each model parameter, where we adopt flat priors for all the parameters. We estimate the best-fit from the median and determine their uncertainties from the 16th and 84th percentile.

Given the spectral resolution of the G140M grating, which is approximately 500 km s^{-1} at the Ly α wavelength, the Ly α lines are mostly unresolved or only marginally resolved. Therefore, we use a Gaussian profile to describe the Ly α line without modeling more complex line structures (such as asymmetries and outflows). This approach is sufficient for measuring the Ly α line fluxes and equivalent widths, but does not allow us to identify any double-peak Ly α emitters, as the peak separation is typically $\lesssim 500 \text{ km s}^{-1}$ (e.g., A. Verhamme et al. 2018; R. P. Naidu et al. 2022).

Note that we aim to test whether galaxies near quasars show less attenuated Ly α lines compared to galaxies outside the quasars' ionized bubble. Therefore, we need to exclude possible AGNs from our sample, as AGNs may also produce their own ionized bubbles and show unattenuated Ly α lines. We identify possible AGNs by measuring the C IV line fluxes (e.g., K. Nakajima et al. 2018; A. Saxena et al. 2022). For each galaxy in our sample, we fit its spectrum at $1500\text{\AA} < \lambda_{\text{rest}} < 1600\text{\AA}$ as a power-law continuum plus a Gaussian emission line. We find five galaxies with rest-frame C IV EW higher than 12\AA ; according to K. Nakajima et al. (2018), these objects are likely AGNs, and we exclude them from the rest of this Section. We present more information about these C IV emitters in the Appendix.

The final sample consists of 50 galaxies at $6 < z < 7$ with $M_{\text{UV}} < -19$, including 28 galaxies in the J0100+2802 field and 22 galaxies in the J1148+5251 field. The median M_{UV} of this sample is -19.97 . Of these, 14 galaxies have a Ly α line detection with $> 3\sigma$ significance. Table 1 summarizes the properties of the galaxies in our sample.

3.2. The Ly α Profile of Galaxies Near Quasars

We now investigate the Ly α line profiles for galaxies near the quasars. Due to the IGM's absorption, most $z > 6$ galaxies have no flux bluer than the systemic Ly α wavelengths. In contrast, for galaxies residing within ionized bubbles generated by luminous quasars, the blue wing of their Ly α lines is expected to be visible due to the low IGM opacity. To investigate this effect, we stack

⁹ <https://jwst-pipeline.readthedocs.io/en/stable/>

¹⁰ <https://github.com/gbrammer/msaexp/tree/main/msaexp>

¹¹ <https://pypeit.readthedocs.io/en/stable/>

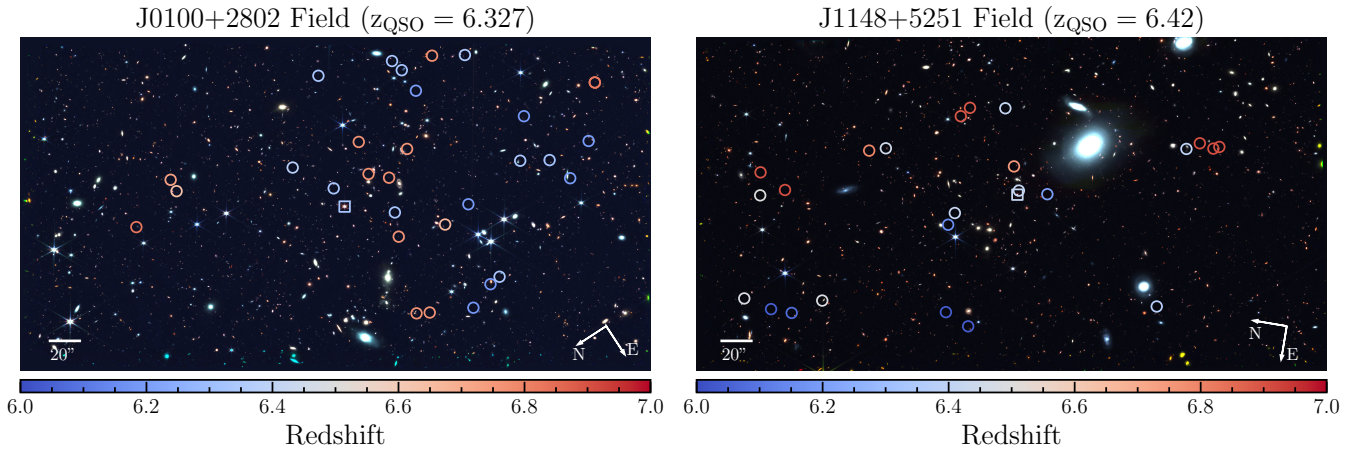


Figure 1. The two quasar fields targeted by this work, J0100+2802 (left) and J1148+5251 (right). We compose the RGB images using the NIRCcam F115W, F200W, and F356W images from the EIGER project (e.g., [D. Kashino et al. 2023](#)). The quasars are at the field centers, marked by the squares. The circles mark the galaxy sample of this study, with the colors representing their systemic redshifts. There are 28 targets and 22 targets in the J0100+2802 and J1148+5251 fields, respectively.

Table 1. Properties of galaxies in our sample

ID	RA (deg)	Dec (deg)	z_{sys}	M_{UV} (mag)	$z_{\text{Ly}\alpha}$	$F_{\text{Ly}\alpha}$ ($10^{-18}\text{erg s}^{-1}\text{cm}^{-2}$)	$\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ (\AA)	$f_{\text{esc}}^{\text{Ly}\alpha}$
J1148.2637	177.112276	52.869027	6.037	-20.45		< 2.76	< 12.19	< 0.03
J1148.3391	177.117208	52.902130	6.075	-20.07		< 0.59	< 2.48	< 0.02
J1148.3461	177.109067	52.873590	6.039	-20.24		< 0.70	< 3.59	< 0.08
J1148.3608	177.117109	52.906070	6.042	-19.19	$6.045^{+0.001}_{-0.001}$	$3.10^{+0.55}_{-0.47}$	$34.34^{+10.88}_{-8.04}$	0.30 ± 0.25
J1148.3822	177.096543	52.834714	6.375	-19.97		< 1.71	< 90.55	< 0.60
J1148.4246	177.111852	52.896855	6.485	-19.63		< 2.50	< 27.07	< 0.27

NOTE—The entire table will be available online upon the acceptance of this paper.

the spectra of galaxies in our sample, as the spectra of individual galaxies do not have sufficient signal-to-noise ratio (S/N). Specifically, we consider three subsets of the galaxies:

1. $|z - z_Q| < \frac{(1+z_Q)\Delta v_{\text{max}}}{c}$, where we take $\Delta v_{\text{max}} = 2500 \text{ km s}^{-1}$ in this work, and z_Q is the quasar’s redshift. This sample is referred to as “galaxies near quasars”;
2. $6 < z < z_Q - \frac{(1+z_Q)\Delta v_{\text{max}}}{c}$, referred to as “foreground galaxies”;
3. $z_Q + \frac{(1+z_Q)\Delta v_{\text{max}}}{c} < z < 7$, referred to as “background galaxies”.

We note here that all the targeted galaxies have transverse distances to the quasars of $r_{\perp} \lesssim 1$ proper Mpc. The expected sizes of the ionized bubble around $z > 6$

luminous quasars is $\gtrsim 1$ Mpc, given the typical quasar lifetimes ($\gtrsim 10^6$ yr; e.g., [A.-C. Eilers et al. 2017](#), also see Eilers et al. submitted). Therefore, we do not apply a cut on r_{\perp} when selecting the “galaxies near quasars” sample.

For each subset, we select galaxies with at least 2σ detection for the Ly α flux. We subtract the best-fit continuum model (from Section 3.1) from the spectrum of each galaxy and shift the spectrum to the rest-frame according to its systemic redshift. We then regrid the spectra to a common wavelength grid and stack the spectra using inverse-variance weighting. The uncertainties of the stacked spectra are derived by propagating the errors of the individual objects’ spectra.

Figure 2 presents the stacked spectra. The stacked spectra of galaxies near quasars show significant flux bluer than the systemic Ly α wavelength. In con-

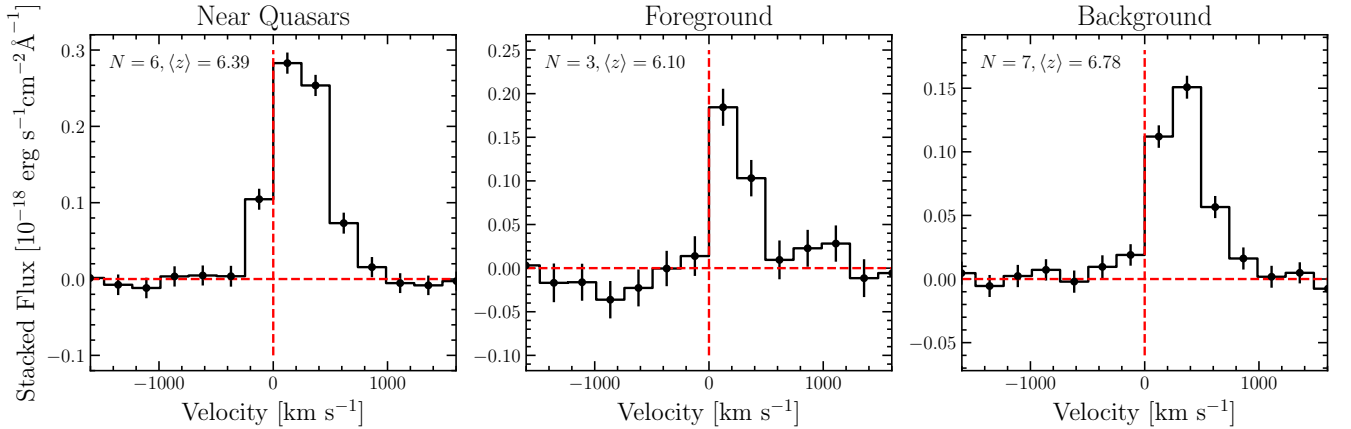


Figure 2. The stacked Ly α line profile for galaxies near quasars (left), foreground galaxies (middle), and background galaxies (right). The text in each panel marks the number of galaxies and the average redshift for the stacked spectra. The dashed lines mark the systemic Ly α wavelength and the level of zero flux. Galaxies near quasars show significant flux bluer than the systemic Ly α wavelength, indicating low IGM opacity around quasars. In contrast, the foreground and background galaxies do not have significant flux bluer than the Ly α wavelength.

trast, the stacked spectra for foreground and background galaxies do not have significant flux blueward of the systemic Ly α wavelength. This result supports the scenario where quasars enhance the Ly α transmission of their surrounding IGM, enabling the detection of Ly α blue wings for galaxies near quasars. Again, we note that the G140M grating does not have sufficient resolution to unveil the complex line structures (e.g., the double peaks) of the Ly α lines.

3.3. The Ly α Equivalent Width Distribution

We then evaluate the rest-frame Ly α equivalent width ($\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$) of the targeted galaxies. If the IGM around quasars has low opacity to Ly α photons, galaxies near high-redshift quasars should show higher $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ compared to other galaxies at similar redshifts.

The top panel of Figure 3 shows the $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ of individual galaxies in our sample. The red dashed line marks the quasars' redshifts, and the shaded region represents the region around the quasars (i.e., $\Delta v < 2500 \text{ km s}^{-1}$). Among the 15 galaxies near quasars, six (40%) of them have Ly α lines detected at higher than 3σ significance. In contrast, this fraction is only $8/35 = 22.9\%$ for foreground and background galaxies. The fractions of galaxies with $\text{EW}_{\text{Ly}\alpha}^{\text{rest}} > 10 \text{ \AA}$ are $5/15 = 33.3\%$ and $7/35 = 20.0\%$ for galaxies near quasars and foreground/background galaxies, respectively. This comparison already hints that galaxies around quasars have less attenuated Ly α emissions, due to the low IGM opacity in the quasar's proximity zone.

To further quantify the difference in Ly α emissions between galaxies inside and outside the quasars' ionized bubbles, we estimate the distribution of $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ for

galaxies in our sample. Following the method in M. Tang et al. (2024b), we assume a log-normal distribution for $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$:

$$p(x|\mu, \sigma) = \frac{A}{\sqrt{2\pi}\sigma x} \times \exp\left[-\frac{(\ln x - \mu)^2}{2\sigma^2}\right] \quad (1)$$

where x is the random variable ($\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ in this case), μ and σ are the mean and standard deviation of the log-normal distribution, and A is the normalization. We derive the posterior distribution of (μ, σ) according to Bayes' theorem:

$$p(\mu, \sigma|\text{obs}) \propto p(\text{obs}|\mu, \sigma)p(\mu, \sigma) \quad (2)$$

where $p(\mu, \sigma)$ is the prior, and we adopt flat priors for μ and σ . $p(\text{obs}|\mu, \sigma)$ is the probability of getting the observed $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ for the galaxies. Specifically,

$$\begin{aligned} p(\text{obs}|\mu, \sigma) &= \prod_i p(\text{obs}, i|\mu, \sigma) \\ &= \prod_i \int_0^{+\infty} p(x)_{\text{obs}, i} \cdot p(x|\mu, \sigma) dx \end{aligned} \quad (3)$$

where i goes through all galaxies in the sample, and $p(x)_{\text{obs}, i}$ is the posterior distribution of the galaxies' $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$. For galaxies with $> 3\sigma$ Ly α line detection, we assume a Gaussian error for the detection, i.e.,

$$p(x)_{\text{obs}, i} = \frac{1}{\sqrt{2\pi}\sigma_{\text{obs}, i}} \cdot \exp\left[-\frac{(x - x_{\text{obs}, i})^2}{2\sigma_{\text{obs}, i}^2}\right] \quad (4)$$

For upper limits, we have

$$p(\text{obs}, i|\mu, \sigma) = p(x < x_{3\sigma, i}^{\text{lim}}|\mu, \sigma) \quad (5)$$

We run Markov Chain Monte Carlo (MCMC) to sample the posterior of (μ, σ) .

Using this method, we model the $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ distributions for galaxies near quasars, foreground galaxies, and background galaxies. For comparison, we also fit the distributions for NIRSPEC-observed $6 < z < 7$ galaxies compiled by [Y. Kageura et al. \(2025\)](#). When fitting the galaxies from [Y. Kageura et al. \(2025\)](#), we apply a cut of $M_{\text{UV}} < -19$ to match the absolute magnitude of our sample.

Table 2 lists the best-fit parameters for these galaxy subsets. Using the posterior of (μ, σ) , we also compute the Ly α -emitter (LAE) fractions (defined as galaxies with $\text{EW}_{\text{Ly}\alpha}^{\text{rest}} > 10\text{\AA}$) for the subsets of galaxies, which are shown in the bottom panel of Figure 3. The foreground and background galaxy samples have LAE fractions close to the $6 < z < 7$ galaxy sample from [Y. Kageura et al. \(2025\)](#). In contrast, in both quasar fields, galaxies near quasars show higher LAE fractions. This comparison again indicates that galaxies near quasars have less-attenuated Ly α lines compared to other $z > 6$ galaxies, due to the low IGM opacity around quasars. Note that the UV slope and UV magnitude distributions of galaxies near quasars are very similar to those of foreground and background galaxies; therefore, the difference in $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ can only be attributed to the difference in IGM opacity, i.e., the influence of the quasars' ionizing radiation fields.

We notice that galaxies near quasars have LAE fractions similar to $z \sim 5$ galaxies measured by [M. Tang et al. \(2024b\)](#). Our result indicates that the unattenuated Ly α line EW of $M_{\text{UV}} < -19$ galaxies has little evolution from $z \sim 6.4$ to $z \sim 5$.

Finally, we note that including C IV emitters into the $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ distribution has a negligible impact on the derived LAE fractions. By including the C IV emitters, the $\text{EW}_{\text{Ly}\alpha}^{\text{rest}} > 10\text{\AA}$ fraction of galaxies near quasars changes to $\chi_{\text{LAE}}(\text{EW} > 10\text{\AA}) = 0.44_{-0.14}^{+0.15}$, and that of foreground + background galaxies changes to $\chi_{\text{LAE}}(\text{EW} > 10\text{\AA}) = 0.24_{-0.06}^{+0.07}$.

4. IMPLICATION FOR THE ESCAPE FRACTION OF REIONIZATION-EPOCH GALAXIES

The enhanced $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ and the detection of the Ly α blue wing for galaxies near quasars indicate that we can detect the Ly α emission line from these galaxies without significant IGM attenuation. These galaxies provide a rare opportunity to directly constrain the Ly α emission lines properties of $z > 6$ galaxies, which is usually impossible due to the high IGM neutral fraction ($x_{\text{HI}} \gtrsim 0.1$, e.g., [X. Jin et al. 2023](#); [D. Āurovčíková et al. 2024](#)). As an example, in this section, we use galaxies near quasars

in our sample to put direct constraints on the Ly α escape fraction ($f_{\text{esc}}^{\text{Ly}\alpha}$) of $z > 6$ galaxies.

To evaluate $f_{\text{esc}}^{\text{Ly}\alpha}$, for each galaxy, we adopt the H β flux from the [O III] emitter catalog presented in [J. Matthee et al. \(2023\)](#) and [D. Kashino et al. \(2025\)](#). We note that one galaxy does not have H β measurement and is excluded from the escape fraction analysis. We obtain the dust attenuation (A_V) from the catalog and correct the H β fluxes for dust attenuation. The median attenuation of the sample is 0.22. We also estimate the continuum flux at $\lambda_{\text{rest}} = 1215.67\text{\AA}$ (denoted by $F_{\text{cont}}^{1215.67\text{\AA}}$) by fitting its F115W and F200W magnitudes using a power law SED. We then compute the escape fraction as

$$f_{\text{esc}}^{\text{Ly}\alpha} = \frac{\text{EW}_{\text{Ly}\alpha}^{\text{rest}} \times F_{\text{cont}}^{1215.67\text{\AA}}}{25 \times F_{\text{H}\beta}} \quad (6)$$

where $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ comes from fitting the MSA spectra, and the coefficient 25 corresponds to case B recombination with temperature $T = 10^4\text{K}$ (e.g., [M. Tang et al. 2024a](#); [Y. Kageura et al. 2025](#)). This approach avoids possible impacts due to the slitloss of the MSA slitlets.

With the $f_{\text{esc}}^{\text{Ly}\alpha}$ of individual galaxies, we compute the distribution of $f_{\text{esc}}^{\text{Ly}\alpha}$ following the same method we use for evaluating the $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ distribution as described in Section 3.3. Specifically, we assume a lognormal distribution for $f_{\text{esc}}^{\text{Ly}\alpha}$, and use Equations 1 to 5 to evaluate the posterior of (μ, σ) for the distribution. For galaxies near quasars, we get $\mu = -2.78_{-1.19}^{+1.48}$ and $\sigma = 3.41_{-1.39}^{+1.06}$. These numbers indicate an average escape fraction of $f_{\text{esc}}^{\text{Ly}\alpha} = 0.14 \pm 0.04$, consistent with the measurements for $z \lesssim 5$ galaxies (e.g., [X. Lin et al. 2024](#)).

The LyC escape fraction is positively correlated with the Ly α escape fraction, and most galaxies have $f_{\text{esc}}^{\text{LyC}} < f_{\text{esc}}^{\text{Ly}\alpha}$ (e.g., [M. Dijkstra et al. 2016](#); [A. Verhamme et al. 2017](#)). We estimate the average LyC escape fraction for galaxies near quasars using the relation in [R. Begley et al. \(2024\)](#), who suggested that $f_{\text{esc}}^{\text{LyC}} = 0.15 \times f_{\text{esc}}^{\text{Ly}\alpha}$. This relation gives $f_{\text{esc}}^{\text{LyC}} \approx 0.02 \pm 0.01$. We note that the $f_{\text{esc}}^{\text{LyC}} - f_{\text{esc}}^{\text{Ly}\alpha}$ relation has large scatter and might depend on the properties of galaxies; therefore, we take $f_{\text{esc}}^{\text{LyC}} < f_{\text{esc}}^{\text{Ly}\alpha} \approx 0.14$ as a conservative upper limit. Due to the small sample size (15 galaxies in total, where 6 galaxies have Ly α detections), we are not able to further explore the correlation between $f_{\text{esc}}^{\text{Ly}\alpha}$ and other galaxy properties (like M_{UV} and UV-slope; e.g., [L. Anderson et al. 2017](#); [J. Chisholm et al. 2022](#)).

Our analysis demonstrates the unique power of high-redshift quasar fields in measuring the escape fraction of EoR galaxies. The estimated $f_{\text{esc}}^{\text{LyC}}$ disfavors the ‘‘high escape fraction’’ scenario with $f_{\text{esc}}^{\text{LyC}} \gtrsim 0.2$ (e.g., [R. P. Naidu et al. 2020](#)), at least for galaxies with $M_{\text{UV}} < -19$

Table 2. The $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ distribution of galaxies

Sample	$\langle z \rangle$	$\mu(\text{EW}_{\text{Ly}\alpha}^{\text{rest}})$ (\AA)	$\sigma(\text{EW}_{\text{Ly}\alpha}^{\text{rest}})$ (\AA)	$\chi_{\text{LAE}}(\text{EW} > 10\text{\AA})$	$\chi_{\text{LAE}}(\text{EW} > 25\text{\AA})$
Galaxies near Quasars	6.37	$2.07^{+0.53}_{-0.89}$	$1.70^{+0.89}_{-0.51}$	$0.45^{+0.14}_{-0.14}$	$0.24^{+0.11}_{-0.08}$
Foreground Galaxies	6.14	$-0.89^{+1.63}_{-1.40}$	$3.14^{+1.10}_{-1.08}$	$0.16^{+0.10}_{-0.06}$	$0.09^{+0.07}_{-0.05}$
Background Galaxies	6.71	$0.03^{+0.91}_{-1.23}$	$2.86^{+1.05}_{-0.78}$	$0.21^{+0.10}_{-0.07}$	$0.12^{+0.08}_{-0.05}$
Foreground + Background	6.53	$0.00^{+0.78}_{-1.20}$	$2.61^{+0.99}_{-0.68}$	$0.19^{+0.06}_{-0.06}$	$0.10^{+0.05}_{-0.04}$
Y. Kageura et al. (2025) ¹	6.42	$0.76^{+0.63}_{-0.92}$	$1.96^{+0.76}_{-0.46}$	$0.22^{+0.07}_{-0.06}$	$0.11^{+0.04}_{-0.04}$

¹Galaxies observed by NIRSpect at $6 < z < 7$, as compiled by Y. Kageura et al. (2025).

NOTE—All samples are selected to have $M_{\text{UV}} < -19$.

at $z \sim 6.4$. Given the UV luminosity functions of $z > 6$ galaxies (e.g., R. J. Bouwens et al. 2022), a low escape fraction of $f_{\text{esc}}^{\text{LyC}} \approx 2\%$ indicates that luminous galaxies (e.g., those with $M_{\text{UV}} < -18$) cannot make a major contribution to the ionizing photon production. In a recent study, C. Simmonds et al. (2024) measured the ionizing photon production efficiency (ξ_{ion}) for $3 < z < 9$ galaxies using NIRCам imaging. With the $\xi_{\text{ion}}(M_{\text{UV}}, z)$ they presented, we only need a low escape fraction $f_{\text{esc}}^{\text{LyC}} \lesssim 0.05$ to reionize the universe. This picture is consistent with our measurement. The estimated $f_{\text{esc}}^{\text{LyC}}$ is also consistent with the finding of C. Papovich et al. (2025), who performed stellar population and nebular emission modeling for $4.5 < z < 9.0$ observed using *JWST* imaging and prism spectroscopy and found an average escape fraction of $f_{\text{esc}}^{\text{LyC}} = 0.03 \pm 0.01$.

4.1. Systematic Uncertainties

In order to measure the $f_{\text{esc}}^{\text{Ly}\alpha}$ for galaxies near quasars, we need to assume that these galaxies have negligible IGM attenuation for their Ly α lines. Here we discuss the systemic uncertainties introduced by this assumption. The Ly α blue wing in the stacked spectra (Figure 2) indicates that the IGM opacity around the quasars is similar to that of $z \sim 5$ IGM (e.g., M. J. Hayes et al. 2021). This similarity is further supported by the recent transverse proximity effect measurement for J0100+2802 (Eilers et al. submitted). In short, Eilers et al. (submitted) measured the transmission flux for background galaxies ($z > z_Q + 0.1$) and found a Ly α transmission of $\mathcal{T} \gtrsim 0.5$ for the IGM near the quasar. In comparison, IGM at $z \lesssim 5$ has transmission of $\mathcal{T} \gtrsim 0.1$ (e.g., R. Thomas et al. 2021; R. A. Meyer et al. 2025). We also notice that the IGM attenuation for Ly α line fluxes is $\lesssim 10\%$ at $z \lesssim 5$ (e.g., M. J. Hayes et al. 2021; J. Matthee et al. 2022). We thus expect the uncertainty introduced by our assumption to be smaller than 10%.

This uncertainty does not influence our main result, i.e., the average Ly α escape fraction is about 0.14.

Another major uncertainty is from the limited sample size, as the ‘‘galaxies near quasars’’ sample only contains 15 galaxies. Future NIRSpect/MSA observations for more high-redshift quasars will reduce the survey variance and provide tighter constraints on the escape fraction of $z > 6$ galaxies.

5. CONCLUSION

We present *JWST* NIRSpect/MSA observations for galaxies in two high-redshift quasar fields, namely J0100+2802 ($z = 6.327$) and J1148+5251 ($z = 6.42$). We analyze 50 galaxies at $6 < z < 7$ with $M_{\text{UV}} < -19$, among which 15 are located near the quasars (with $\Delta v < 2500 \text{ km s}^{-1}$). We measure the Ly α fluxes, profiles, and rest-frame equivalent widths of these galaxies. Leveraging the low IGM opacity around quasars, we are able to measure the Ly α lines from galaxies near quasars without significant IGM attenuation. This allows us to put direct constraints on the Ly α escape fraction of $z > 6$ galaxies. Our main findings include

1. The stacked Ly α spectrum for galaxies near quasars clearly shows a significant emission bluer than the systemic Ly α wavelength, while foreground and background galaxies show no significant emission bluer than the systemic Ly α wavelength in their stacked spectra. This result indicates that the Ly α lines from galaxies near high-redshift quasars are less attenuated compared to other galaxies at similar redshifts.
2. The LAE fraction of galaxies near quasars is $\chi_{\text{LAE}}(\text{EW} > 10\text{\AA}) = 0.45^{+0.14}_{-0.14}$, higher than that of foreground and background galaxies ($0.19^{+0.06}_{-0.06}$). This result confirms that galaxies near quasars

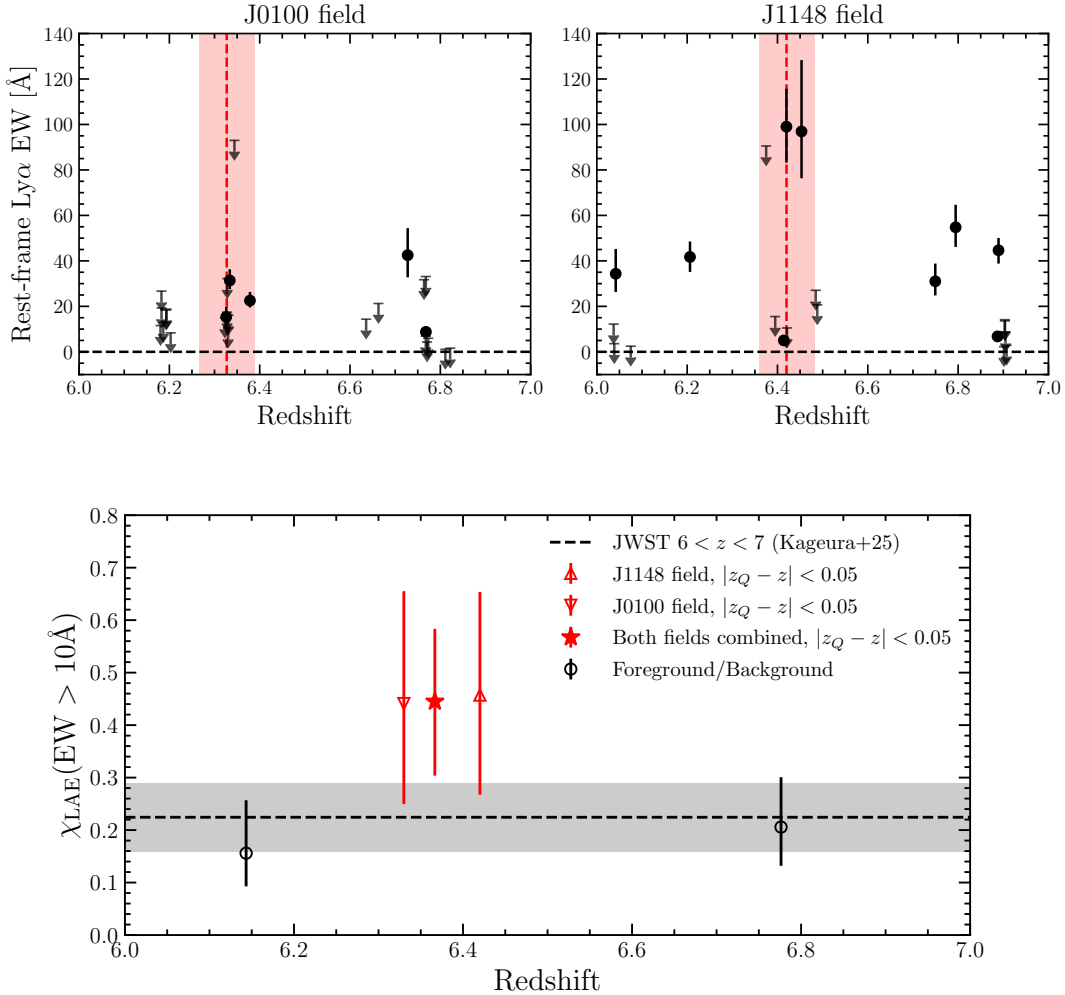


Figure 3. *Top Panels:* redshift and $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ distribution of galaxies in the J0100 field and the J1148 field, in the left and right panel, respectively. The red shaded region marks the quasar’s vicinity $|\Delta v| < 2500 \text{ km s}^{-1}$. *Bottom:* The fraction of galaxies with $\text{EW}_{\text{Ly}\alpha}^{\text{rest}} > 10\text{Å}$. This fraction is estimated by fitting the galaxy $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ with a log-normal distribution (see text for details). We also include the NIRSpec-observed galaxies at $6 < z < 7$ from Y. Kageura et al. (2025) for comparison. In both quasar fields, galaxies in the quasar’s vicinity show higher LAE fractions than foreground and background galaxies and the general galaxy population. This comparison indicates that IGM around luminous quasars has high transmission to Ly α photons.

have less attenuated Ly α lines. The $\text{EW}_{\text{Ly}\alpha}^{\text{rest}}$ distribution of galaxies near quasars is similar to $z \sim 5$ galaxies.

3. We estimate a median Ly α escape fraction of $f_{\text{esc}}^{\text{Ly}\alpha} = 0.14 \pm 0.04$ for galaxies near quasars. Adopting the scaling relation of $f_{\text{esc}}^{\text{Ly}^{\text{C}}} = 0.15 \times f_{\text{esc}}^{\text{Ly}\alpha}$ from R. Begley et al. (2024), we estimate $f_{\text{esc}}^{\text{Ly}^{\text{C}}} = 0.02 \pm 0.01$ for galaxies near quasars. This result favors the low escape fraction scenario ($f_{\text{esc}}^{\text{Ly}^{\text{C}}} \lesssim 0.1$) for $z > 6$ galaxies.

The ionized bubbles surrounding luminous quasars serve as unique laboratories for studying galaxy evolution and reionization in the early universe. High-resolution spectroscopy with NIRSpec’s high-resolution grating or future extremely large telescopes will be able

to resolve the Ly α line profiles of galaxies near high-redshift quasars (e.g., R. P. Naidu et al. 2022). Such observations will enable key measurements for EoR galaxies, such as outflow properties and interstellar medium conditions. These experiments will offer valuable insights into the physical properties of EoR galaxies.

The *JWST* data used in this paper can be found in MAST: [10.17909/wm0b-5n06](https://mast.stsci.edu/portal/#/home/jwst/10.17909/wm0b-5n06). All the code and data used in this work will be available online upon the acceptance of this paper.

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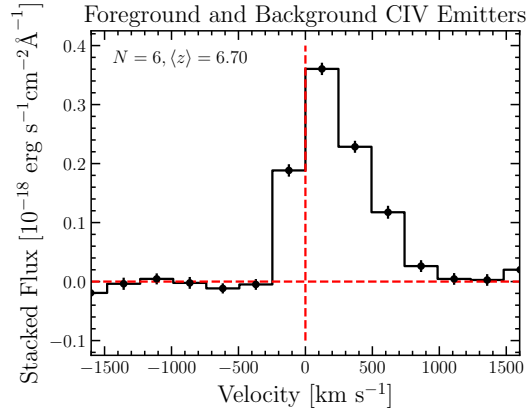


Figure 4. The stacked Ly α line for C IV emitters at the quasars’ foreground and background, showing significant flux bluer than the systemic Ly α wavelength. This result indicate that C IV emitters can produce their own ionized bubbles.

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Facilities: JWST(NIRCam, NIRSpect)

Software: PyeIt, msaexp, jwst, astropy, bilby(G. Ashton et al. 2019)

APPENDIX

A. C IV EMITTERS IN OUR SAMPLE

In addition to the 50 galaxies in the main sample, we also identified five C IV emitters with $EW_0(\text{C IV}) > 12\text{\AA}$. These galaxies are possible AGNs. Three C IV emitters show Ly α emissions at $> 3\sigma$ levels, two of which have $EW_{\text{Ly}\alpha}^{\text{rest}} > 100\text{\AA}$. We highlight object J0100_10287 at $z = 6.77$, which has $EW_{\text{Ly}\alpha}^{\text{rest}} = 152_{-30}^{+60}\text{\AA}$ and is the strongest Ly α emitter in the MASQUERADE sample with $M_{\text{UV}} < -19$. The properties of the C IV emitters are also included in Table 1.

Following the method in Section 3.2, we produce the stacked Ly α profile for the C IV emitters at the quasars’ foreground and background. Although these galaxies are outside the quasars’ ionized bubbles, the stacked Ly α spectrum shows a significant Ly α blue wing. This result indicates that these C IV emitters can produce their own ionized bubbles. Our finding is consistent with previous studies that suggest C IV emitters might be strong Ly α leakers (e.g., R. P. Naidu et al. 2022; D. Schaerer et al. 2022; S. Mascia et al. 2023).

We also note that only one C IV emitter is located in the ionized bubbles around quasars. Including this object in the “galaxies near quasars” sample has a negligible impact on the estimated $EW_{\text{Ly}\alpha}^{\text{rest}}$ and $f_{\text{esc}}^{\text{Ly}\alpha}$ distributions.

REFERENCES

- Anderson, L., Governato, F., Karcher, M., Quinn, T., & Wadsley, J. 2017, MNRAS, 468, 4077, doi: 10.1093/mnras/stx709
- Ashton, G., et al. 2019, Astrophys. J. Suppl., 241, 27, doi: 10.3847/1538-4365/ab06fc
- Atek, H., Labbé, I., Furtak, L. J., et al. 2024, Nature, 626, 975, doi: 10.1038/s41586-024-07043-6
- Begley, R., Cullen, F., McLure, R. J., et al. 2024, MNRAS, 527, 4040, doi: 10.1093/mnras/stad3417
- Bosman, S. E. I., Kakiichi, K., Meyer, R. A., et al. 2020, ApJ, 896, 49, doi: 10.3847/1538-4357/ab85cd
- Bosman, S. E. I., Davies, F. B., Becker, G. D., et al. 2022, MNRAS, 514, 55, doi: 10.1093/mnras/stac1046

- Bouwens, R. J., Illingworth, G., Ellis, R. S., Oesch, P., & Stefanon, M. 2022, *ApJ*, 940, 55, doi: [10.3847/1538-4357/ac86d1](https://doi.org/10.3847/1538-4357/ac86d1)
- Bouwens, R. J., Illingworth, G. D., Oesch, P. A., et al. 2012, *ApJL*, 752, L5, doi: [10.1088/2041-8205/752/1/L5](https://doi.org/10.1088/2041-8205/752/1/L5)
- Brammer, G. 2023,, 0.6.17 Zenodo, doi: [10.5281/zenodo.8319596](https://doi.org/10.5281/zenodo.8319596)
- Chen, Z., Stark, D. P., Mason, C. A., et al. 2025, arXiv e-prints, arXiv:2505.24080, doi: [10.48550/arXiv.2505.24080](https://doi.org/10.48550/arXiv.2505.24080)
- Chisholm, J., Saldana-Lopez, A., Flury, S., et al. 2022, *MNRAS*, 517, 5104, doi: [10.1093/mnras/stac2874](https://doi.org/10.1093/mnras/stac2874)
- Davies, F. B. 2020, *MNRAS*, 494, 2937, doi: [10.1093/mnras/staa528](https://doi.org/10.1093/mnras/staa528)
- Dijkstra, M., Gronke, M., & Venkatesan, A. 2016, *ApJ*, 828, 71, doi: [10.3847/0004-637X/828/2/71](https://doi.org/10.3847/0004-637X/828/2/71)
- Eilers, A.-C., Davies, F. B., Hennawi, J. F., et al. 2017, *ApJ*, 840, 24, doi: [10.3847/1538-4357/aa6c60](https://doi.org/10.3847/1538-4357/aa6c60)
- Fan, X., Strauss, M. A., Becker, R. H., et al. 2006, *AJ*, 132, 117, doi: [10.1086/504836](https://doi.org/10.1086/504836)
- Finkelstein, S. L., D'Aloisio, A., Paardekooper, J.-P., et al. 2019, *ApJ*, 879, 36, doi: [10.3847/1538-4357/ab1ea8](https://doi.org/10.3847/1538-4357/ab1ea8)
- Hayes, M. J., Runnholm, A., Gronke, M., & Scarlata, C. 2021, *ApJ*, 908, 36, doi: [10.3847/1538-4357/abd246](https://doi.org/10.3847/1538-4357/abd246)
- Izotov, Y. I., Chisholm, J., Worseck, G., et al. 2022, *MNRAS*, 515, 2864, doi: [10.1093/mnras/stac1899](https://doi.org/10.1093/mnras/stac1899)
- Izotov, Y. I., Worseck, G., Schaerer, D., et al. 2018, *MNRAS*, 478, 4851, doi: [10.1093/mnras/sty1378](https://doi.org/10.1093/mnras/sty1378)
- Jaskot, A. E., Silveyra, A. C., Plantinga, A., et al. 2024, *ApJ*, 973, 111, doi: [10.3847/1538-4357/ad5557](https://doi.org/10.3847/1538-4357/ad5557)
- Jin, X., Yang, J., Fan, X., et al. 2023, *ApJ*, 942, 59, doi: [10.3847/1538-4357/aca678](https://doi.org/10.3847/1538-4357/aca678)
- Kageura, Y., Ouchi, M., Nakane, M., et al. 2025, arXiv e-prints, arXiv:2501.05834, doi: [10.48550/arXiv.2501.05834](https://doi.org/10.48550/arXiv.2501.05834)
- Kashino, D., Lilly, S. J., Matthee, J., et al. 2023, *ApJ*, 950, 66, doi: [10.3847/1538-4357/acc588](https://doi.org/10.3847/1538-4357/acc588)
- Kashino, D., Lilly, S. J., Matthee, J., et al. 2025, arXiv e-prints, arXiv:2506.03121, doi: [10.48550/arXiv.2506.03121](https://doi.org/10.48550/arXiv.2506.03121)
- Lin, X., Cai, Z., Wu, Y., et al. 2024, *ApJS*, 272, 33, doi: [10.3847/1538-4365/ad3e7d](https://doi.org/10.3847/1538-4365/ad3e7d)
- Lu, T.-Y., Mason, C. A., Mesinger, A., et al. 2024, arXiv e-prints, arXiv:2411.04176, doi: [10.48550/arXiv.2411.04176](https://doi.org/10.48550/arXiv.2411.04176)
- Ma, X., Hopkins, P. F., Kasen, D., et al. 2016, *MNRAS*, 459, 3614, doi: [10.1093/mnras/stw941](https://doi.org/10.1093/mnras/stw941)
- Ma, X., Kasen, D., Hopkins, P. F., et al. 2015, *MNRAS*, 453, 960, doi: [10.1093/mnras/stv1679](https://doi.org/10.1093/mnras/stv1679)
- Mascia, S., Pentericci, L., Calabrò, A., et al. 2023, *A&A*, 672, A155, doi: [10.1051/0004-6361/202345866](https://doi.org/10.1051/0004-6361/202345866)
- Mason, C. A., Treu, T., Dijkstra, M., et al. 2018, *ApJ*, 856, 2, doi: [10.3847/1538-4357/aab0a7](https://doi.org/10.3847/1538-4357/aab0a7)
- Mason, C. A., Fontana, A., Treu, T., et al. 2019, *MNRAS*, 485, 3947, doi: [10.1093/mnras/stz632](https://doi.org/10.1093/mnras/stz632)
- Matthee, J., Mackenzie, R., Simcoe, R. A., et al. 2023, *ApJ*, 950, 67, doi: [10.3847/1538-4357/acc846](https://doi.org/10.3847/1538-4357/acc846)
- Matthee, J., Sobral, D., Gronke, M., et al. 2018, *A&A*, 619, A136, doi: [10.1051/0004-6361/201833528](https://doi.org/10.1051/0004-6361/201833528)
- Matthee, J., Naidu, R. P., Pezzulli, G., et al. 2022, *MNRAS*, 512, 5960, doi: [10.1093/mnras/stac801](https://doi.org/10.1093/mnras/stac801)
- Meyer, R. A., Roberts-Borsani, G., Oesch, P., & Ellis, R. S. 2025, arXiv e-prints, arXiv:2504.02683, doi: [10.48550/arXiv.2504.02683](https://doi.org/10.48550/arXiv.2504.02683)
- Naidu, R. P., Forrest, B., Oesch, P. A., Tran, K.-V. H., & Holden, B. P. 2018, *MNRAS*, 478, 791, doi: [10.1093/mnras/sty961](https://doi.org/10.1093/mnras/sty961)
- Naidu, R. P., Tacchella, S., Mason, C. A., et al. 2020, *ApJ*, 892, 109, doi: [10.3847/1538-4357/ab7cc9](https://doi.org/10.3847/1538-4357/ab7cc9)
- Naidu, R. P., Matthee, J., Oesch, P. A., et al. 2022, *MNRAS*, 510, 4582, doi: [10.1093/mnras/stab3601](https://doi.org/10.1093/mnras/stab3601)
- Nakajima, K., Schaerer, D., Le Fèvre, O., et al. 2018, *A&A*, 612, A94, doi: [10.1051/0004-6361/201731935](https://doi.org/10.1051/0004-6361/201731935)
- Nikolić, I., Mesinger, A., Mason, C. A., et al. 2025, arXiv e-prints, arXiv:2501.07980, doi: [10.48550/arXiv.2501.07980](https://doi.org/10.48550/arXiv.2501.07980)
- Oke, J. B., & Gunn, J. E. 1983, *ApJ*, 266, 713, doi: [10.1086/160817](https://doi.org/10.1086/160817)
- Papovich, C., Cole, J. W., Hu, W., et al. 2025, arXiv e-prints, arXiv:2505.08870, doi: [10.48550/arXiv.2505.08870](https://doi.org/10.48550/arXiv.2505.08870)
- Planck Collaboration, Aghanim, N., Akrami, Y., et al. 2020, *A&A*, 641, A6, doi: [10.1051/0004-6361/201833910](https://doi.org/10.1051/0004-6361/201833910)
- Prochaska, J. X., Hennawi, J. F., Westfall, K. B., et al. 2020, *Journal of Open Source Software*, 5, 2308, doi: [10.21105/joss.02308](https://doi.org/10.21105/joss.02308)
- Protušová, K., Bosman, S. E. I., Wang, F., et al. 2024, arXiv e-prints, arXiv:2412.12256, doi: [10.48550/arXiv.2412.12256](https://doi.org/10.48550/arXiv.2412.12256)
- Robertson, B. E. 2022, *ARA&A*, 60, 121, doi: [10.1146/annurev-astro-120221-044656](https://doi.org/10.1146/annurev-astro-120221-044656)
- Rosdahl, J., Blaizot, J., Katz, H., et al. 2022, *MNRAS*, 515, 2386, doi: [10.1093/mnras/stac1942](https://doi.org/10.1093/mnras/stac1942)
- Saxena, A., Cryer, E., Ellis, R. S., et al. 2022, *MNRAS*, 517, 1098, doi: [10.1093/mnras/stac2742](https://doi.org/10.1093/mnras/stac2742)
- Schaerer, D., Izotov, Y. I., Worseck, G., et al. 2022, *A&A*, 658, L11, doi: [10.1051/0004-6361/202243149](https://doi.org/10.1051/0004-6361/202243149)
- Sharma, M., Theuns, T., Frenk, C., et al. 2016, *MNRAS*, 458, L94, doi: [10.1093/mnras/slw021](https://doi.org/10.1093/mnras/slw021)

- Simmonds, C., Tacchella, S., Hainline, K., et al. 2024, MNRAS, 535, 2998, doi: [10.1093/mnras/stae2537](https://doi.org/10.1093/mnras/stae2537)
- Steidel, C. C., Bogosavljević, M., Shapley, A. E., et al. 2018, ApJ, 869, 123, doi: [10.3847/1538-4357/aaed28](https://doi.org/10.3847/1538-4357/aaed28)
- Tang, M., Stark, D. P., Topping, M. W., Mason, C., & Ellis, R. S. 2024a, ApJ, 975, 208, doi: [10.3847/1538-4357/ad7eb7](https://doi.org/10.3847/1538-4357/ad7eb7)
- Tang, M., Stark, D. P., Ellis, R. S., et al. 2024b, MNRAS, 531, 2701, doi: [10.1093/mnras/stae1338](https://doi.org/10.1093/mnras/stae1338)
- Thomas, R., Pentericci, L., Le Fèvre, O., et al. 2021, A&A, 650, A63, doi: [10.1051/0004-6361/202038438](https://doi.org/10.1051/0004-6361/202038438)
- Torralba, A., Matthee, J., Naidu, R. P., et al. 2024, A&A, 689, A44, doi: [10.1051/0004-6361/202450318](https://doi.org/10.1051/0004-6361/202450318)
- Umeda, H., Ouchi, M., Kageura, Y., et al. 2025, arXiv e-prints, arXiv:2504.04683, doi: [10.48550/arXiv.2504.04683](https://doi.org/10.48550/arXiv.2504.04683)
- Đurovčíková, D., Eilers, A.-C., Chen, H., et al. 2024, ApJ, 969, 162, doi: [10.3847/1538-4357/ad4888](https://doi.org/10.3847/1538-4357/ad4888)
- Verhamme, A., Orlitová, I., Schaerer, D., et al. 2017, A&A, 597, A13, doi: [10.1051/0004-6361/201629264](https://doi.org/10.1051/0004-6361/201629264)
- Verhamme, A., Garel, T., Ventou, E., et al. 2018, MNRAS, 478, L60, doi: [10.1093/mnras/sly058](https://doi.org/10.1093/mnras/sly058)
- Whitler, L., Stark, D. P., Endsley, R., et al. 2024, MNRAS, 529, 855, doi: [10.1093/mnras/stae516](https://doi.org/10.1093/mnras/stae516)
- Wold, I. G. B., Malhotra, S., Rhoads, J., et al. 2022, ApJ, 927, 36, doi: [10.3847/1538-4357/ac4997](https://doi.org/10.3847/1538-4357/ac4997)
- Yang, J., Wang, F., Fan, X., et al. 2020, ApJ, 904, 26, doi: [10.3847/1538-4357/abbc1b](https://doi.org/10.3847/1538-4357/abbc1b)